

1 **Diagnosis of Relative Humidity Changes in a Warmer Climate**

2 **Using Tracers of Last Saturation**

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## ABSTRACT

6 The zonal mean relative humidity response to a doubling of CO<sub>2</sub> in a climate model is  
7 examined, using two different methods to differentiate the effects of circulation changes from  
8 spatially inhomogeneous temperature changes. The tropical and subtropical response are  
9 found to be largely dependent on circulation changes, particularly a poleward expansion  
10 and deepening of the Hadley circulation, a poleward shift of the extratropical jets, and an  
11 increase in the height of the tropopause. The responses near the extratropical tropopause and  
12 in the lower troposphere are largely dependent on changes in the distribution and gradients  
13 of temperature.

# 14 1. Introduction

15 Climate models indicate that the water vapor feedback is roughly equivalent to that  
16 expected from constant global mean relative humidity (RH) (Soden and Held 2006; Randall  
17 et al. 2007). Analyses of observed climate variations in the recent historical record are  
18 consistent with this view (Soden et al. 2002; Dessler et al. 2008). A constant global mean  
19 RH does not necessarily correspond to a static distribution of RH, however, and even small  
20 changes can be consequential for other aspects of the climate (e.g., Sherwood et al. 2009).

21 Relative humidity is an important factor in determining the distribution and occurrence  
22 of clouds (Sundqvist 1978; Price and Wood 2002). An increase in the fraction of optically  
23 thin high clouds with warming would represent a positive feedback, as such clouds are rel-  
24 atively transmissive to sunlight, largely opaque to outgoing longwave radiation, and have  
25 a substantially different emission temperature than the surface. The converse is true for  
26 low clouds, as the increase in solar albedo that they provide outweighs their effect as long-  
27 wave absorbers (Manabe and Strickler 1964; Hartmann et al. 1992; Chen et al. 2000). A  
28 greater understanding of the underlying causes of simulated RH changes and their plausi-  
29 bility may therefore be helpful in constraining cloud feedbacks, which currently represent  
30 the largest source of inter-model spread in climate sensitivity (Randall et al. 2007). Inho-  
31 mogeneity in RH changes also impacts the distribution of both latent and radiative heating,  
32 which may then project onto the atmospheric circulation (Schneider et al. 2009), large-scale  
33 convective organization (Gray and Jacobson 1977), and the level at which deep convection  
34 detrains (Hartmann and Larson 2002). Regional shifts of the climatological distribution of  
35 RH thus have the potential to influence climate on a wide range of scales.

36 Several studies have noted that the simulated RH response to warming exhibits a distinc-  
37 tive zonal mean pattern (Mitchell and Ingram 1992; Lorenz and DeWeaver 2007a; Sherwood  
38 et al. 2009). This pattern is characterized by a horseshoe-shaped decrease of relative humidity  
39 throughout the tropical upper troposphere, subtropics, and extratropical free troposphere,  
40 with a slightly more pronounced decrease in the southern hemisphere. This horseshoe is  
41 bracketed by an increase of relative humidity in the tropical mid-troposphere and extratrop-  
42 ical tropopause layer, as shown in Fig. 1. The RH difference shown in Fig. 1 is averaged over  
43 ten models in the World Climate Research Program’s (WCRP’s) Coupled Model Intercom-  
44 parison Project phase 3 (CMIP3) multi-model dataset, and is calculated using time mean  
45 zonal mean relative humidities from the final five years of the slab ocean control (SlabCNTL)  
46 and doubled carbon dioxide (2xCO<sub>2</sub>) experiments. Although the details of the signal vary  
47 among constituent models, the qualitative pattern is largely robust.

48 Relative humidity is defined in this analysis as the ratio of specific humidity to satu-  
49 ration specific humidity. Saturation specific humidity is a function of local temperature as  
50 expressed by the Clausius-Clapeyron equation. Free tropospheric specific humidity is in turn  
51 determined to leading order by the large-scale circulation and temperature fields, with con-  
52 densate evaporation playing a minor role (Sherwood 1996; Pierrehumbert and Roca 1998;  
53 Galewsky et al. 2005; Sherwood and Meyer 2006; Wright et al. 2009b).

54 Unsaturated air parcels conserve specific humidity. To the extent that condensate evapo-  
55 ration is unimportant, the specific humidity in unsaturated air equals the saturation specific  
56 humidity at the point at which it was last saturated. Galewsky et al. (2005, hereafter GSH)  
57 used “tracers of last saturation” to trace the pathways taken by air parcels since their points  
58 of last saturation. This technique provides insight into the ways in which the circulation and

59 temperature fields together determine the distributions of atmospheric specific and relative  
60 humidity. Here we apply an updated formulation of this tracer technique to the output of  
61 two integrations of a GCM, one of which simulates modern climate and one of which sim-  
62 ulates a climate with doubled CO<sub>2</sub>. The results help to establish the relative influences of  
63 shifts in atmospheric circulation as compared to inhomogeneous changes in temperature on  
64 the characteristic RH response shown in Fig. 1.

## 65 **2. Model Configuration**

66 This study employs a two-level global modeling procedure to investigate the mechanisms  
67 behind relative humidity change in a warmer climate. First, a GCM is run to provide six-  
68 hourly circulation and temperature fields that are representative of both a modern (CTL) and  
69 doubled CO<sub>2</sub> (WRM) climate. These meteorology fields are then used as inputs to a global  
70 tracer transport model, which includes an independent hydrologic cycle and is outfitted with  
71 a last saturation tracer scheme (GSH; Hurley and Galewsky 2009).

### 72 *a. General Circulation Model*

73 The base meteorology for this study is generated by two integrations of the Goddard In-  
74 stitute for Space Studies (GISS) ModelE (Schmidt et al. 2006). The first of these, designated  
75 CTL, uses atmosphere and ocean conditions consistent with the year 1979, including green-  
76 house gas concentrations and sea surface temperatures. The second simulation is designated  
77 WRM, and is performed using a slab ocean version of the model with atmospheric CO<sub>2</sub>

78 doubled from the 1979 value at the outset. The concentration of atmospheric CO<sub>2</sub> is held  
79 constant at 337.1 ppmv for the CTL simulation and 674.2 ppmv for the WRM simulation.

80 Both model simulations are run at 2° latitude by 2.5° longitude resolution with 20 vertical  
81 levels. Advection of temperature and water vapor conserves potential enthalpy and mass, and  
82 is computed using a quadratic upstream scheme with nine higher-order moments (Prather  
83 1986). This yields an effective tracer resolution of approximately 0.7°×0.8°. The model  
84 physics and radiation are described in detail by Schmidt et al. (2006).

85 Sea surface temperatures and sea ice extent for the CTL simulation follow a fixed annual  
86 cycle averaged over 1975 to 1984, with all other boundary conditions set to 1979 values as  
87 discussed by Schmidt et al. (2006). Atmospheric CO<sub>2</sub> is also kept constant at 1979 levels.  
88 This simulation is run for ten years; output from the last five years is used for this analysis.

89 The WRM simulation is similar to the CTL simulation, with the addition of a mixed layer  
90 heat flux model (q-flux) and doubled CO<sub>2</sub>. Initial mixed layer heat transport is prescribed  
91 using implied values from a five-year climatology generated during the CTL simulation. Net  
92 global heating at the surface during the CTL run is 0.09 W m<sup>-2</sup>, well within the ±0.5 W m<sup>-2</sup>  
93 threshold recommended for a q-flux setup run (Schmidt et al. 2006). The mixed layer  
94 depth varies according to a fixed seasonal cycle and is assumed to be isothermal. Energy  
95 is conserved by incorporating fluxes between the mixed layer and a deeper layer between  
96 the base of the current mixed layer and the base of the mixed layer at its annual maximum  
97 depth. Sea surface temperatures and sea ice extent are determined dynamically during the  
98 model integration.

99 Instantaneous meteorological variables are saved every six hours during each model run.  
100 Saved surface variables include orography, surface geopotential, surface temperature, surface

101 pressure, latent heat flux, sensible heat flux, and the zonal and meridional components of sur-  
102 face stress. Atmospheric variables are saved at all 20 vertical levels and include temperature,  
103 specific humidity, and zonal and meridional winds. For compatibility with the tracer trans-  
104 port model, these data are interpolated from the ModelE's  $2^\circ \times 2.5^\circ$  latitude-longitude grid  
105 to a T42 Gaussian grid using bilinear interpolation. The vertical coordinate is unchanged.

106 *b. Tracer Transport Model*

107 Tracer transport is accomplished using the offline Model for Atmospheric Transport  
108 and Chemistry (MATCH) developed at the National Center for Atmospheric Research  
109 (NCAR) (Rasch et al. 1997). The MATCH model uses a semi-Lagrangian advection scheme,  
110 and includes an independent hydrologic cycle with parameterizations for cloud physics and  
111 convection.

112 The MATCH integrations presented here are performed using a 30 minute timestep, with  
113 linear interpolation between the six-hourly meteorological fields. The model is run on a T42  
114 Gaussian horizontal grid with 20 hybrid sigma vertical levels, matching the input meteoro-  
115 logical data. Tracer advection is calculated using a semi-Lagrangian transport scheme with  
116 enforced mass conservation (Rasch and Williamson 1990; Rasch et al. 1995). Subgrid-scale  
117 turbulent mixing is represented by a vertical eddy diffusion parameterization.

118 The parameterizations for clouds and convection are based on those developed for ver-  
119 sion 3 of the NCAR Community Climate Model (CCM3). In particular, MATCH uses the  
120 prognostic cloud parameterization presented by Rasch and Kristjansson (1998) and the con-  
121 vection scheme described by Hack et al. (1998). The convective parameterization partitions

122 convective transport into deep convection (Zhang and McFarlane 1995) and shallow convec-  
 123 tion (Hack 1994). Tracers are advected both within the convective parameterization and by  
 124 the large-scale circulation.

125 *c. Tracer Formulation*

126 The base formulation of the tracers of last saturation follows that of GSH. Specifically, a  
 127 set of  $N$  zonally symmetric tracer domains is chosen to cover the global troposphere. Each  
 128 grid point is associated with the domain that contains it; we will call the tracer associated  
 129 with this domain the local tracer ( $\mathcal{L}$ ) and all others nonlocal tracers ( $\mathcal{T}_i; i = 1, \dots, N - 1$ ).

130 All tracers are initially set to zero. During model integration, whenever free tropospheric  
 131 RH exceeds a saturation threshold of 90% the local tracer is set to one and all nonlocal  
 132 tracers are set to zero at that point:

$$\begin{aligned}\mathcal{L}(\lambda, \phi, p, t) &= 1 \\ \mathcal{T}_i(\lambda, \phi, p, t) &= 0|_{i=1, \dots, N-1} \\ \mathcal{S}(\lambda, \phi, p, t) &= 0 \\ \mathcal{S}_{\text{amt}}(\lambda, \phi, p, t) &= 0,\end{aligned}$$

133 where  $\lambda$  and  $\phi$  represent the longitude and latitude of the saturated grid cell,  $p$  represents  
 134 the vertical coordinate, and  $t$  denotes the model timestep.  $\mathcal{S}$  and  $\mathcal{S}_{\text{amt}}$  are the surface source  
 135 tracers, which are defined below. Whenever the local RH is below the threshold value, the  
 136 tracers of last saturation are permitted to advect and mix unchanged. A RH threshold of



137 90% is chosen to reflect the fact that saturation operates at spatial scales well below the  
 138 grid scale; that is, some air parcel in the grid volume may be at saturation even though the  
 139 mean RH for the entire volume is below 100%. The results are insensitive within reasonable  
 140 perturbations to this threshold ( $\pm 10\%$  RH). Saturation is determined according to MATCH's  
 141 internal hydrologic cycle, rather than the GCM output.

142 The evaporative source at the surface is incorporated by treating the lowest model layer  
 143 separately: all last saturation tracers in this layer are set to zero and a source tracer ( $\mathcal{S}$ ) is  
 144 defined with a value equal to the current specific humidity ( $q$ ) in the grid cell:

$$\begin{aligned}\mathcal{T}_i(\lambda, \phi, p_b, t) &= 0|_{i=1, \dots, N} \\ \mathcal{S}(\lambda, \phi, p_b, t) &= q(\lambda, \phi, p_b, t) \\ \mathcal{S}_{\text{amt}}(\lambda, \phi, p_b, t) &= 1,\end{aligned}$$

145  $p_b$  denotes the lowest model layer. Note that the local tracer  $\mathcal{L}$  is replaced by  $\mathcal{S}$  at the  
 146 surface, so that there are  $N$  nonlocal tracers rather than  $N - 1$ . For bookkeeping purposes  
 147 we also define a source amount tracer  $\mathcal{S}_{\text{amt}}$  that follows the definitions of  $\mathcal{L}$  and  $\mathcal{T}$ . The  
 148 source tracer is permitted to mix, so that the value of  $\mathcal{S}$  at any location may reflect several  
 149 excursions to the surface.

150 The local specific humidity  $q(\lambda, \phi, p, t)$  can then be reconstructed via the linear combi-  
 151 nation

$$q(\lambda, \phi, p, t) = \mathcal{L}(\lambda, \phi, p, t)q^*(\lambda, \phi, p, t) + \sum_i^{N-1} \mathcal{T}_i(\lambda, \phi, p, t)\langle q_i^* \rangle + \mathcal{S}(\phi, p) \quad (1)$$

152 where  $q^*(\lambda, \phi, p, t)$  is the local saturation mixing ratio and  $\langle q_i^* \rangle$  represents the density weighted

153 mean saturation specific humidity for tracer domain  $i$ . The tracers generally obey the con-  
 154 straint

$$\mathcal{L}(\lambda, \phi, p, t) + \sum_i^{N-1} \mathcal{T}_i(\lambda, \phi, p, t) + \mathcal{S}_{\text{amt}}(\lambda, \phi, p, t) = 1 \quad (2)$$

155 in our simulations; after a brief initial spin-up period significant deviations from the occur  
 156 only in the stratosphere. RH is then reconstructed as

$$\text{RH}(\lambda, \phi, p, t) = \frac{q(\lambda, \phi, p, t)}{q^*(\lambda, \phi, p, t)} \quad (3)$$

157 with  $q(\lambda, \phi, p, t)$  determined by Eq. 1. GSH discuss technical issues involved in this recon-  
 158 struction and quantify several sources of error.

159 Figure 2 shows a direct comparison between the modeled and reconstructed zonal mean  
 160 RH fields for the MATCH integration using CTL meteorological fields as input. The quali-  
 161 tative patterns match up remarkably well, and the point to point comparison also indicates  
 162 excellent agreement in both the tropics and extratropics. Excluding the model layers be-  
 163 low 900 hPa and above 110 hPa, where boundary layer or stratospheric influences render  
 164 the reconstruction less effective, the Pearson correlation coefficients between modeled and  
 165 reconstructed RH are greater than 0.95.

166 Figure 2 includes two adjustments to the tracer scheme presented by GSH. First, we  
 167 have altered the distribution of tracer domains (defined by dotted black lines). Although the  
 168 chosen domains remain zonally axisymmetric, they now provide global coverage (as opposed  
 169 to 50°S to 50°N in GSH). The horizontal resolution of the tracer domains is approximately  
 170 5° latitude equatorward of 50°, with a  $\sim 15^\circ$  domain out to 65° and a 25° domain extending

171 to the pole in both hemispheres. We also increase the tracer domain resolution with altitude,  
 172 so that the vertical domain sizes are roughly equivalent in  $\log(p)$  space. This allows us to  
 173 better diagnose the mechanisms influencing upper tropospheric humidity, particularly in the  
 174 tropics. Second, we have applied a temperature correction to the online tracer calculation.  
 175 The GSH formulation predicted extremely high humidities in the upper troposphere. This  
 176 bias resulted from the transport of trace amounts of source and lower tropospheric tracer  
 177 into the upper troposphere. Although these tracer concentrations were quite small, they  
 178 were associated with values of  $\langle q^* \rangle$  that were comparatively quite high, and thus exerted a  
 179 disproportionately large influence on the reconstructed humidity. We have addressed this  
 180 issue by including an online calculation of density weighted mean temperature  $\langle T_i \rangle$  for each  
 181 tracer domain. At each timestep, if  $\langle T_i \rangle > T(\lambda, \phi, p)$ , then that tracer is converted to local  
 182 tracer:

$$\mathcal{L}(\lambda, \phi, p) = \mathcal{L}(\lambda, \phi, p) + \mathcal{T}_i(\lambda, \phi, p)$$

$$\mathcal{T}_i(\lambda, \phi, p) = 0|_{\langle T_i \rangle > T(\lambda, \phi, p)}.$$

183 This adjustment compensates for any spurious vertical tracer transport that does not lead  
 184 to grid scale saturation, and is physically equivalent to assuming that condensate is imme-  
 185 diately removed at the subgrid scale in both parameterized convective updrafts and vertical  
 186 advection.

### 3. Simulated Climate Changes

Investigation of the mechanisms underlying the characteristic pattern shown in the multimodel mean RH changes (Fig. 1) requires meteorological output from a GCM that exhibits this pattern of change as climate warms. Figure 3 shows the RH difference between the WRM and CTL runs of the GISS ModelE. The changes between these two runs agrees quite well with the multimodel mean change, both qualitatively and quantitatively. This indicates that the ModelE is a reasonable choice for examining the root causes behind the pattern of RH changes.

A number of studies have examined the distribution of circulation and temperature changes in the model simulations submitted to the CMIP3 intercomparison project (e.g., Randall et al. 2007; Lorenz and DeWeaver 2007b; Vecchi and Soden 2007; Lu et al. 2008; Gastineau et al. 2008). In general, these studies find that the tropopause height increases, the tropical overturning circulation expands poleward, deepens, and weakens, the subtropical jets shift poleward, the lapse rate of temperature with altitude is reduced, and the equator-to-pole temperature gradient in the upper troposphere is increased. These findings are generally supported by observational studies that focus on recent trends in atmospheric temperature and circulation (e.g., Santer et al. 2003; Seidel and Randel 2006; Hu and Fu 2007), although there is some disagreement with observed trends in the strength of the Hadley cell (Mitas and Clement 2006).

Figure 4 shows simulated circulation and temperature changes between the CTL and WRM runs of the ModelE GCM. The troposphere warms nearly everywhere (Fig. 4a; shading), with the strongest warming in the tropical upper troposphere. This reduces the tropo-

209 spheric lapse rate in the tropics and subtropics. The equator-to-pole temperature gradient  
210 is reduced in the lower troposphere and increased in the upper troposphere. Changes in  
211 zonal mean streamfunction (Fig. 4a; black contours) indicate that the Hadley cell expands  
212 and deepens. Taken together with the warming in the tropical upper troposphere, the lat-  
213 ter is at least qualitatively consistent with the fixed anvil temperature hypothesis, which  
214 postulates that tropical convective detrainment is constrained to occur at roughly the same  
215 temperature as climate changes (Hartmann and Larson 2002). The strength of the Hadley  
216 circulation is very similar between the CTL and WRM slab ocean simulations of the ModelE,  
217 with a slight strengthening or weakening depending on the metric used. Vecchi and Soden  
218 (2007) and others report that the tropical overturning circulation weakens in the CMIP3  
219 models; however, this decrease is primarily manifested in the Walker circulation rather than  
220 the Hadley cell. The strength of the longitudinal Walker circulation in the ModelE decreases  
221 in WRM simulation (not shown), consistent with this consensus. Gastineau et al. (2008)  
222 show that changes in the strength of the Hadley cell are much more variable in CMIP3  
223 models, and Mitas and Clement (2006) report recent positive trends in reanalysis data that  
224 are not reproduced by GCM simulations of twentieth century climate. In the context of  
225 current scientific understanding, the representation of changes in the strength of the tropical  
226 overturning circulation is reasonable and consistent with expectations.

227 Shifts in the zonal mean zonal wind (Fig. 4a; white contours) indicate that the subtrop-  
228 ical jets intensify and shift poleward in the ModelE, consistent with the CMIP3 multimodel  
229 mean (Lorenz and DeWeaver 2007b). Figure 4b shows zonal mean tropopause height and  
230 temperature changes for the World Meteorological Organization (WMO; 1957) tropopause.  
231 Tropopause pressure (height) decreases (increases) globally with a minimum shift in the trop-

232 ics, while tropopause temperature increases everywhere but in the deep tropics, which exhibit  
233 a slight cooling. Both of these are consistent with the the CMIP3 multimodel means (Lorenz  
234 and DeWeaver 2007b) and with observations of recent anthropogenic trends (Santer et al.  
235 2003).

236 If the multimodel means and existing observational studies are considered as a baseline  
237 consensus, the circulation and temperature changes simulated by the GISS ModelE are  
238 generally consistent with this consensus. One important caveat is that ozone levels in the  
239 simulations presented here are fixed at 1979 values, so the impacts of stratospheric ozone  
240 recovery on circulation and humidity changes (Son et al. 2008, 2009) are not included in this  
241 analysis.

## 242 4. Tracer Model Experiments

243 The GCM simulations provide a means by which to describe the control and doubled  
244 CO<sub>2</sub> climates. The MATCH tracer transport model is employed at the second level because  
245 it affords greater flexibility. By separating calculations of the tracer distribution from the  
246 circulation and temperature fields that determine them, the mechanisms that control the  
247 distribution of RH changes can be better identified and isolated. It is therefore important  
248 that the results of the MATCH runs using GISS output are similarly able to reproduce the  
249 expected pattern of RH change. Figure 5 shows the zonal mean difference of RH between  
250 the WRM and CTL runs for the MATCH model hydrologic cycle. The contents of this figure  
251 closely mirror those of Fig. 3.

252 The MATCH runs are forced using ModelE output. This leaves two significant differences

253 between the simulations used to prepare Fig. 3 and those used to prepare Fig. 5: the advec-  
254 tion scheme and the hydrologic cycle parameterizations. The close correspondence between  
255 the pattern of RH change in MATCH and that in the GCM suggests that the distribution  
256 of RH changes is not strongly sensitive to the details of these two parameterizations. This  
257 conclusion is supported by the robust nature of the pattern among the CMIP3 model simu-  
258 lations, which also contain a variety of advection and microphysical parameterizations. We  
259 note, however, that the MATCH model and most of the CMIP3 models indicate a stronger  
260 RH signal throughout the troposphere than the ModelE. Humidity in the ModelE exhibits  
261 a fairly strong dependence on condensate evaporation (Wright et al. 2009b), together with  
262 much higher ice water paths than either observations or other GCMs (Waliser et al. 2009).  
263 The occurrence of condensate evaporation is strongly dependent on ambient RH: less con-  
264 densate evaporates at higher RH while more condensate evaporates at lower RH (Dessler and  
265 Sherwood 2004; Wright et al. 2009a). It is not possible to differentiate the signals sufficiently  
266 to draw firm conclusions, but the evidence suggests that the weaker signal in the ModelE  
267 may be attributed to a greater role of condensate evaporation in that GCM. It follows then  
268 that if condensate evaporation exerts a larger influence on humidity than is currently be-  
269 lieved then the strength of the RH signal shown in Fig. 1 will be suppressed. This would  
270 nudge the climate closer to true constant RH, likely leading to a slight increase the strength  
271 of the water vapor feedback (Soden et al. 2002; Minschwaner and Dessler 2004; Minschwaner  
272 et al. 2006).

273 Relative humidity at any given point will remain constant under climate change so long  
274 as the saturation mixing ratio changes by the same fraction at the point in question and at  
275 the relevant point(s) of last saturation. Loosely speaking, this would occur over the entire

276 atmosphere if the circulation remains relatively constant and the temperature changes are  
277 spatially uniform. On the other hand, the circulation does change and temperature changes  
278 have spatial structure in climate model simulations of warming (e.g., Randall et al. 2007;  
279 Lorenz and DeWeaver 2007b; Lu et al. 2008). We wish to attribute the changes in RH shown  
280 in Fig. 1 to these two factors. To what extent are these changes driven by circulation shifts,  
281 and to what extent are they driven by spatially inhomogeneous temperature changes?

282 As a brute force method of separating the roles of circulation and temperature, we run  
283 the MATCH model with temperature and circulation fields chosen from different GCM sim-  
284 ulations. MATCH is run with WRM temperatures and CTL circulation and vice versa. This  
285 is dynamically inconsistent, since temperature and winds are related through the equations  
286 of motion. On the other hand, it is kinematically acceptable for the purpose of diagnosing  
287 the mechanisms controlling water vapor; the water vapor simply evolves in space and time  
288 according to a given set of temperature and wind fields. This approach leverages the offline  
289 tracer transport to separate temperature and circulation in a way that could not be done in  
290 a dynamically consistent calculation.

291 Figure 6a shows RH changes between the CTL MATCH simulation and a simulation in  
292 which WRM atmospheric temperatures are combined with the CTL circulation. This differ-  
293 ence does not show the characteristic horseshoe-shaped pattern of RH decrease throughout  
294 the troposphere, particularly in the tropical and subtropical upper troposphere. It does  
295 capture the RH increases near the extratropical tropopause and in the lower stratosphere,  
296 however, and many aspects of the lower tropospheric response. Figure 6b shows the same  
297 quantity for a MATCH simulation in which the WRM circulation is combined with CTL at-  
298 mospheric temperatures. In this case, the tropical and subtropical free tropospheric response



299 is captured quite well, although the RH decrease near the extratropical tropopause does not  
300 appear. Figure 6 thus indicates that RH changes in the tropical and subtropical troposphere  
301 are dominated by circulation changes, whereas the increase near the extratropical tropopause  
302 and changes in near-surface RH are controlled by temperature changes.

303 The high latitude tropospheric response is far too strong in both perturbation simulations;  
304 in fact, Fig. 3 and Fig. 5 indicate that changes in these regions should be small and of  
305 variable sign, rather than the strongly negative response shown in both panels of Fig. 6.  
306 This mismatch may be a consequence of the shallow convective parameterization in MATCH,  
307 which can moisten the troposphere without causing grid-scale saturation and which is not  
308 handled explicitly by our tracer technique. It could also simply be a result of nonlinearities:  
309 tracer transport is linear, but saturation is nonlinear. Thus, the results in Figs. 6a and 6b  
310 need not add up to those in Fig. 5 in general, although they do so (approximately) in the  
311 tropics and subtropics, and at the extratropical tropopause.

## 312 5. Last Saturation Tracers

313 Figure 2 indicates that the tracer reconstruction of RH agrees both qualitatively and  
314 quantitatively with the RH simulated by MATCH using CTL meteorology. This agreement  
315 translates to RH changes between the CTL and WRM runs, as shown in Fig. 7. The  
316 reconstruction captures much of the structure observed in both the GISS model (Fig. 3)  
317 and the online MATCH hydrologic cycle (Fig. 5), in particular the horseshoe-shaped RH  
318 decrease and the increases in the tropical middle troposphere and extratropical tropopause  
319 layer. The tracer reconstruction of RH does not capture the increase of RH in the lower

320 stratosphere, but this is unsurprising since our choice of tracer domains effectively omits the  
321 stratosphere.

322 This agreement provides a check on the consistency of the tracer formulation. Since the  
323 RH reconstruction successfully captures the pattern of RH changes in the warmer climate,  
324 the tracers can be applied to diagnose some of the relevant mechanisms.

325 A simple diagnostic that can be constructed from the last saturation tracers involves  
326 separating contribution of changes in the local tracer  $\mathcal{L}(\lambda, \phi, p)$ , which represents the amount  
327 of air in a grid cell  $(\lambda, \phi, p)$  that was last saturated with the tracer domain containing that  
328 cell, from that of all nonlocal tracers  $\sum_i^{N-1} \mathcal{T}_i(\lambda, \phi, p)$ . Although it only makes use of a  
329 small fraction of the information carried by the tracers, this diagnostic appears to explain  
330 a large portion of the RH change. Figure 8 shows the zonal and time mean change in the  
331 concentration of  $\mathcal{L}$ . The pattern of changes in  $\mathcal{L}$  agrees remarkably well with the pattern of  
332 changes in simulated RH. In particular, if the proportion of air that is last saturated locally  
333 decreases then the RH tends to decrease, and vice versa. This correspondence is expected,  
334 as air that was last saturated nearby is likely to be closer to saturation now.

335 The close correspondence between the pattern of changes in RH and the pattern of  
336 changes in  $\mathcal{L}$  is particularly relevant near the extratropical tropopause in both hemispheres.  
337 These regions experience an increase in the concentration of local tracer  $\mathcal{L}$ , which acts to  
338 increase RH by increasing the contribution of  $q^*(\phi, p)$  to  $q(\phi, p)$  (Eqn. 1). This is driven in  
339 large part by the gradient of temperature changes in the upper troposphere. In the CTL  
340 simulation, humidity near the extratropical tropopause is determined to a significant extent  
341 by equatorward zones of last saturation. The greatest warming occurs in the tropical upper  
342 troposphere, with diminished warming toward the poles (Fig. 4). This gradient of warming

343 acts to increase the local control of humidity near the extratropical and polar tropopause;  
344 this increase of local control results in an increase of RH. The increase of RH in this region  
345 may in turn lead to an increase in the occurrence of high thin clouds near the extratropical  
346 tropopause, with implications for cloud radiative forcing. Examination of the GCM results  
347 indicates that total cloud cover increases in these regions by 1% to 5% (not shown), due to  
348 an increase in the occurrence of cirrus ice clouds that are formed in situ.

349 As with the MATCH perturbation simulations presented in Section 4, the tracer recon-  
350 struction of RH (Eqs. 1 and 3) can be broken down into two components: one representing the  
351 circulation (the tracers), and one representing temperatures ( $q^*$ ) (cf. Hurley and Galewsky  
352 2009). This attribution is not clean, because the temperature field influences the tracers as  
353 well; if the circulation were held fixed, changes in temperature would change the locations at  
354 which saturation occurs, thus changing the tracer fields. Nonetheless, the correspondence of  
355 many aspects of the results below with those in the previous section — in which an entirely  
356 different method with different limitations was used to separate the roles of temperature and  
357 circulation — suggests that there is some validity to the conclusions.

358 Figure 9 shows zonal mean changes in reconstructed RH using WRM calculations of  $q^*$   
359 with CTL tracers (Fig. 9a), and WRM tracers with CTL  $q^*$  (Fig. 9b).

360 These results support the conclusions drawn from Fig. 6. In particular, circulation  
361 changes appear to play a dominant role in RH changes in the tropical troposphere, while  
362 inhomogeneous changes in temperature appear to control the RH increase near the extratrop-  
363 ical tropopause. These responses can be illustrated in further detail using the last saturation  
364 tracer distributions.

365 Figure 10 shows changes in the concentration of two sets of tracers in the tropical upper

366 troposphere. The first set (left panels) is associated with the layer between 288 hPa and  
367 212 hPa, while the second set (right panels) is associated with the layer above (212 hPa to  
368 150 hPa). There is a dramatic transfer of influence from the lower level to the upper one;  
369 the concentration of the lower set of tracers decreases throughout the tropical troposphere  
370 and appears to be largely replaced by tracer from the upper set. This transfer represents  
371 an upward shift in the zones of last saturation throughout the tropics, consistent with an  
372 upward shift in the tropopause as shown in Fig. 4b, and as expected from the fixed anvil  
373 temperature hypothesis (Hartmann and Larson 2002).

374 Figure 11 shows changes in tracer concentrations associated with humidity in the sub-  
375 tropical free troposphere. This figure shows that the primary regions of last saturation for  
376 the subtropical dry zones shift upward and poleward in the warmer climate. These changes  
377 are consistent with an expansion of the tropical Hadley cell and a poleward shift in the  
378 jetstreams, as shown in Figure 4. In particular, both circulation and tracer shifts are more  
379 pronounced in the southern hemisphere. The decrease of RH in the southern hemisphere is  
380 also stronger in the ModelE, MATCH, tracer reconstruction, and CMIP3 multimodel mean.  
381 As mentioned above, these model runs do not include stratospheric ozone recovery, so these  
382 asymmetries are likely due to differences in the distribution of the continents and orography  
383 between the northern and southern hemispheres.

## 384 6. Summary

385 The zonal mean signature of the relative humidity response to a doubling of CO<sub>2</sub> is  
386 qualitatively robust across climate models. This signature is characterized by a horseshoe-

387 shaped decrease of relative humidity in the tropical upper troposphere, subtropics, and  
388 extratropical free troposphere, with a stronger decrease in the southern hemisphere, and an  
389 increase of RH in the tropical mid-troposphere and extratropical tropopause layer.

390 Two climate model simulations are performed, one of modern climate and one with  
391 doubled CO<sub>2</sub>. Humidity and circulation changes between these simulations are generally  
392 representative of the model simulations submitted to the CMIP3 model intercomparison  
393 project. Six-hourly meteorological output from the GCM simulations is used to drive a three-  
394 dimensional offline tracer transport model that contains both an independent hydrologic  
395 cycle and a zonally axisymmetric last saturation tracer scheme. The tracers are capable of  
396 quantitatively and qualitatively capturing both the modeled RH field and the pattern of RH  
397 response to warming. Two different methods are then used to separate the role of circulation  
398 from that of temperature.

399 Two perturbation simulations are performed using the tracer transport model that pair  
400 modern circulation with doubled CO<sub>2</sub> temperatures and vice versa. The results of these  
401 simulations indicate that the horseshoe-shaped pattern of RH decrease is driven primarily  
402 by circulation shifts, particularly in the tropical and subtropical upper troposphere, while  
403 RH increases near the extratropical tropopause and changes near the surface appear to be  
404 controlled by inhomogeneities in the temperature response to a doubling of CO<sub>2</sub>. Similar  
405 conclusions are reached by manipulating the tracer reconstruction of RH to better differenti-  
406 ate between the contributions of circulation, local temperature, and nonlocal temperatures.

407 Much of the zonal mean RH response is captured by the binary distinction between local  
408 and nonlocal last saturation tracers; that is, if the amount of air in a grid cell that was  
409 last saturated nearby increases, the RH generally increases as well, and vice versa. This

410 correspondence is particularly relevant near the extratropical tropopause, which exhibits  
411 an increase in RH that is associated primarily with an increase in local last saturation.  
412 Both of these are driven in large part by the gradient of temperature changes in the upper  
413 troposphere and at the tropopause, and lead to an increase in high clouds with substantial  
414 implications for cloud radiative forcing in the extratropics and polar regions.

415 The last saturation tracers are used to illustrate the influence of simulated circulation  
416 shifts on zonal mean RH. In particular, last saturation zones for the tropical upper tropo-  
417 sphere shift upward in the doubled CO<sub>2</sub> climate, resulting in a RH decrease. This shift is  
418 consistent with the upward shift of the tropopause and the deepening of tropical convec-  
419 tion associated with the Hadley Cell observed in the simulation. Similarly, the tracers of  
420 last saturation that control RH in the subtropical dry zones shift upward and poleward in  
421 the warmer climate, consistent with a poleward expansion of the tropical circulation and a  
422 poleward shift of the extratropical jets.

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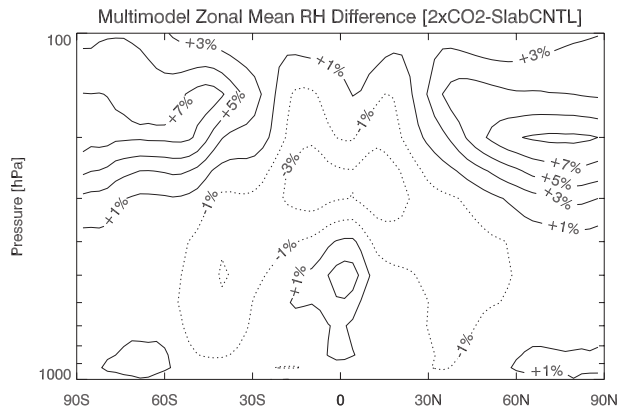


FIG. 1. Zonal mean changes in RH averaged over ten different slab ocean GCMs. Dotted contours represent decreases in the  $2\times\text{CO}_2$  runs as compared to the control runs. Contour intervals are 2% RH; the first dashed contour represents a 1% absolute decrease in RH and the first solid contour represents a 1% absolute increase in RH.



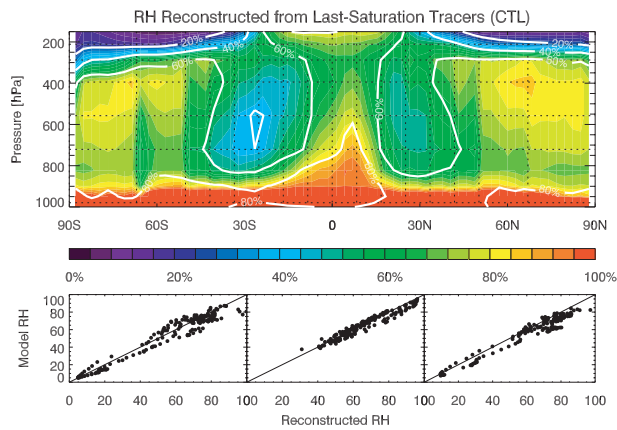


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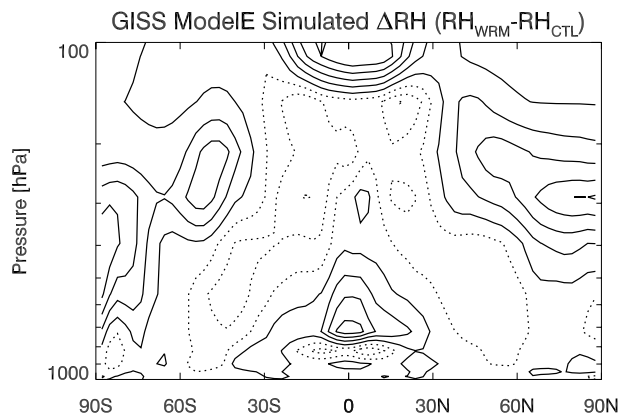


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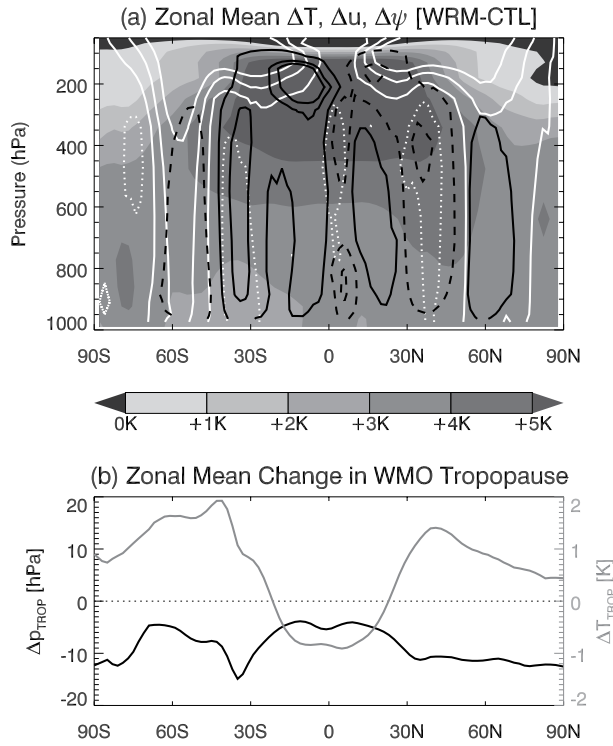


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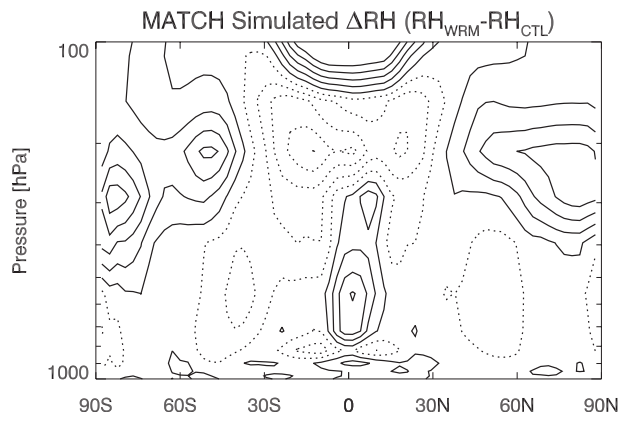


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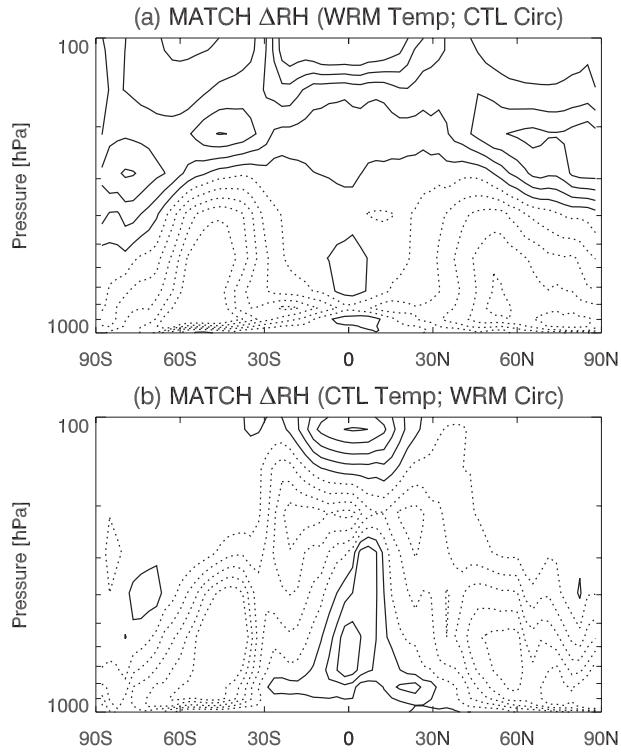


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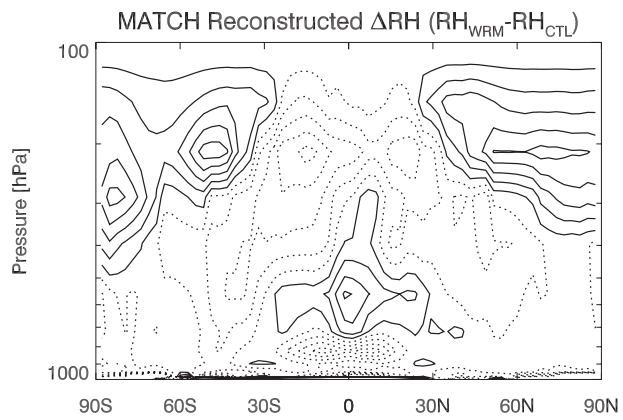


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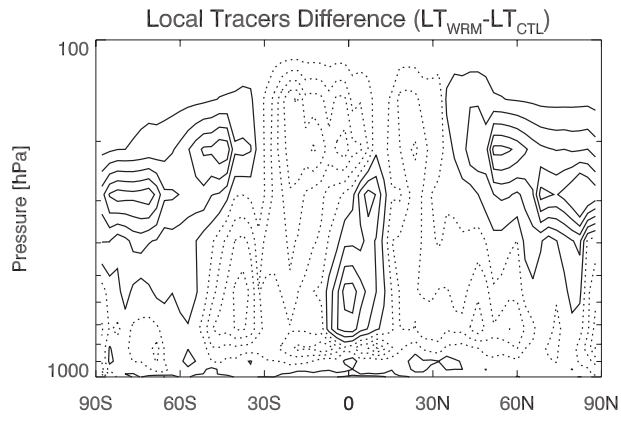


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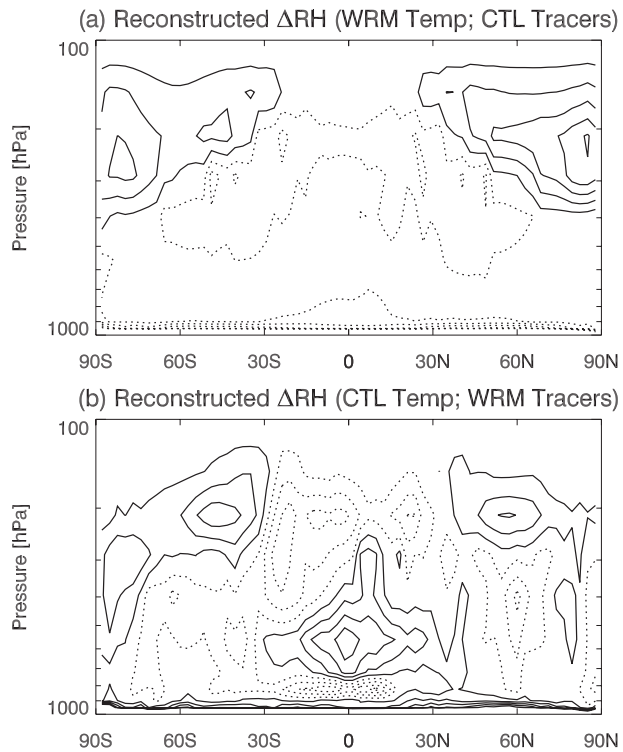


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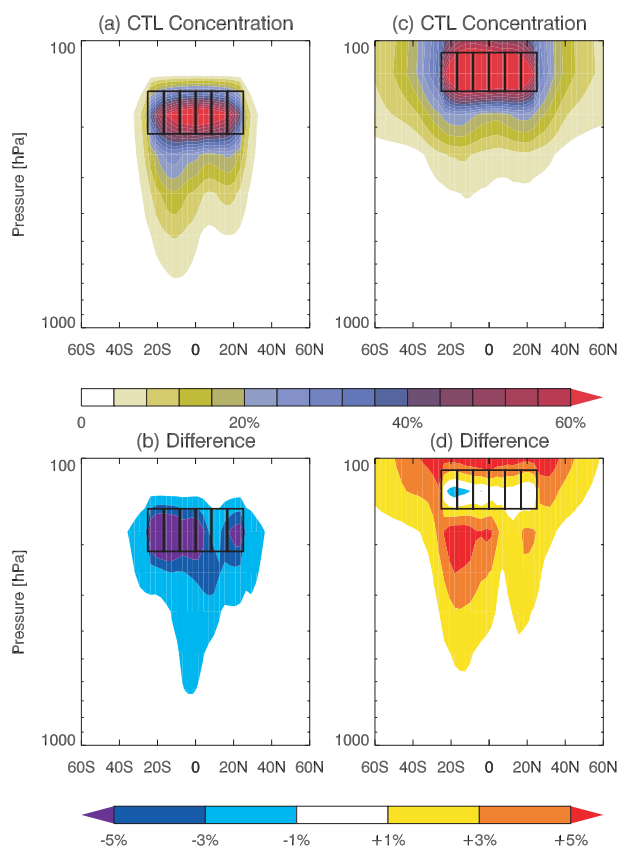


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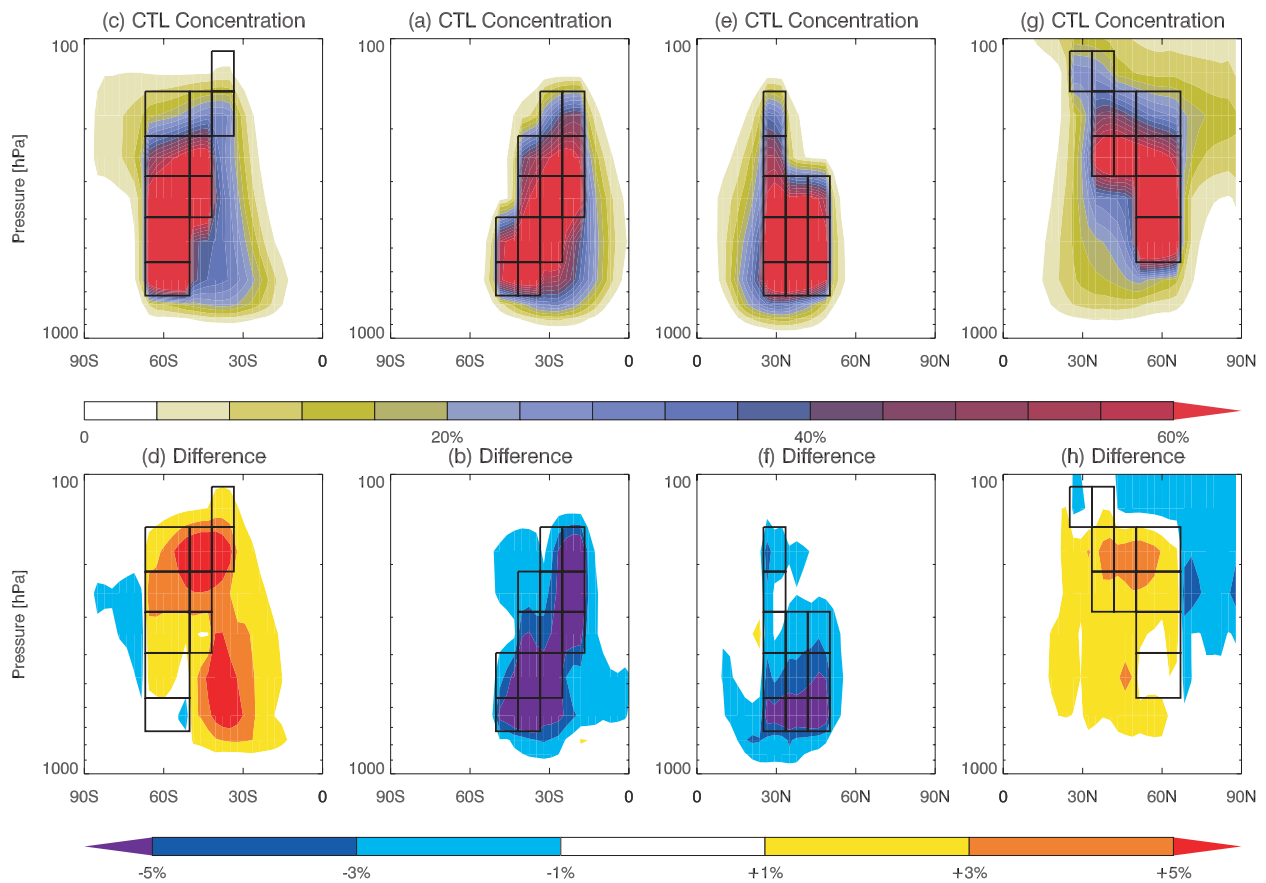


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