Kinematic Constraints on Glacier Contributions to 21st-Century Sea-Level Rise

W. T. Pfeffer,¹,² J. T. Harper,² S. O’Neel²

On the basis of climate modeling and analogies with past conditions, the potential for multimter increases in sea level by the end of the 21st century has been proposed. We consider glaciological constraints on sea-level rise (SLR) during the late Little Ice Age, notably the Maunder Minimum (2, 27). The contrast between strongly reduced SST in the western basin and much less reduced SST in the central Mediterranean basin during the LGM (Fig. 1) indicates a preferentially meridional geographic circulation, with a polar trough that frequently protruded into the western Mediterranean. Such a circulation is favored by northward extension of the Azores High toward Iceland (North Atlantic ridge) or Greenland, blocking moisture supply by the westerlies. It is further enhanced by expansion and intensification of the Siberian High in winter during glacial times (26). A similar configuration is thought to have been common during the late Little Ice Age, notably the Maunder Minimum (2, 27). The invasion of polar air as shown by our data, channeled by the topography of mountain ranges and ice sheets in Europe, would have generated cyclone formation in the Gulf of Genoa more frequently than at present, enhancing precipitation along various storm tracks in easterly directions. Our observations do not support a straightforward zonal LGM atmospheric circulation, as inferred from climate models (19, 28). Instead, we propose that frequent meridional circulation during cold seasons (characterized by the LGM ELA pattern) may have alternated with more zonal circulation during warm seasons. A more comprehensive quantitative assessment of the preferential LGM atmospheric circulation requires the use of both nested model simulation and high-resolution global climate model studies (4, 5, 8, 28) which should fully resolve the changing topography of glaciated mountain ranges and ice sheets. The validation of such models with our three-dimensional LGM climate proxy data ranging from the sea surface to alpine altitudes is a great future challenge.

References and Notes
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Supporting Online Material
www.sciencemag.org/cgi/content/full/321/5763/DC1
Data Supplements S1 to S3
Supporting Online Material

Eustatic ice contributions to sea-level change come from surface mass balance (SMB) losses and discharge of ice into the ocean through marine-terminating glaciers. Dynamically forced discharge, via fast flow and calving of marine-terminating glaciers allowing rapid land-to-ocean transfer of ice, is well known from studies of temperate marine-terminating glaciers (1–4) and is observed in Greenland (5–7). The consensus estimate of sea-level rise (SLR) by 2100 (0.18 to 0.6 m) that was published in the Intergovernmental Panel on Climate Change (IPCC) Fourth Assessment (8) excluded dynamic effects on the grounds that present understanding of the relevant processes is too limited for reliable model estimates. Because modeling (9) and paleo-climate comparisons (10) have yielded multimeter per century estimates of SLR, similar increases have been inferred as a viable 21st-century scenario. Also argued is that feedbacks unaccounted for in the IPCC estimate could quickly cause several meters of very rapid SLR (11, 12).

Accurate SLR forecasts on the century timescale are imperative for planning constructive and cost-effective responses. Underestimates will prompt inadequate preparation for change, whereas overestimates will exhaust and redirect resources inappropriately. Raising California Central Valley levees only 0.15 m, for example, will cost over $1 billion (13); the nonlinearly increasing costs of raising levees 2 m or more without clear and compelling cause would entail enormous expenditures otherwise used for different responses as demanded by a smaller but still significant SLR. We address the plausibility of very rapid SLR from land ice occurring this century. We give
particular emphasis to Greenland because of its vulnerability to ongoing Arctic warming and meltwater-related feedbacks, recent accelerations of ice motion, and its large volume reductions during the last interglacial (14). By using a simple kinematic approach, we determined Greenland and Antarctic outlet glacier velocities required to achieve various magnitudes of SLR by 2100.

To begin, we postulated sea-level increases of 2 and 5 m by 2100 forced solely by Greenland. The total water mass losses required to achieve these targets are 7.24 × 10^7 Gt and 1.81 × 10^8 Gt, respectively (Table 1). Of this mass, 2.58 × 10^6 Gt (less than 4% of the total for 2 m SLR) will be lost as SMB by 2100, which we estimated by integrating present-day values of mass loss forward at present-day rates of change (15), with present-day SMB estimated at 30% of total present-day rate of mass loss (5). Because future SMB is highly uncertain, we also scaled total SMB losses up by a factor of 10 to investigate the effect of uncertainty in SMB. Adjusting total mass losses for SMB contributions yields the mass to be discharged through marine-based outlets (Table 1). Even when scaled up by an order of magnitude, SMB is a very small fraction of the total loss required to produce the targeted SLR. Thus, even large uncertainties in future SMB have little influence on this calculation.

Rapid, dynamically unstable discharge of ice through calving is restricted to glaciers with beds based below sea level. We identified and calculated the aggregate cross-sectional area of Greenland’s marine-terminating outlet glaciers by using surface and bed topography (16) and measured ice velocities (5) to identify all potential pathways for rapid discharge, including channels presently flowing rapidly as well as potentially unstable channels (Fig. 1 and table S1). Cross-sectional areas (gates) for each outlet were calculated at the point of greatest lateral constriction by bedrock in the glacier’s marine-based reach. Ice stream widths in Antarctica can vary in time, but for Greenland outlet glaciers cross-sectional areas are constrained almost entirely by bedrock topography. Of the 290 km^2 total aggregate gate cross-sectional area, we identified 170 km^2 as the aggregate marine-based gate area where drainage to the ocean is not blocked by near-coastal sills standing above present-day sea level. All dynamic discharge (Table 2) must pass through these gates by 2100 to meet 2- and 5-m SLR targets. We considered four scenarios: velocities were calculated for both the “marine-based” gate (170 km^2) and the “total aggregate” gate (290 km^2) given both projected SMB and 10× inflated SMB losses. We then considered whether those velocities are realistic.

Coarseness of the digital elevation models (DEMs) used for surface and bed topography (16) led to uncertainties in the calculated gate areas, which may be substantial but cannot be evaluated directly. We accounted for a potential underestimate of gate area with a calculation using the total aggregate gate area. The total aggregate area exceeds the more relevant marine-based gate area by 70%. Uncertainties arising from the DEM more likely conceal small unresolved channels than large ones, so the actual gate areas may be smaller than we calculate (thus yielding higher velocities).

Table 1. Fluxes and discharges for SLR targets. Q indicates total discharge to 2100 (Gt); q, Q converted to ice flux rate, total to 2100 (km^3/year); Q_2, total dynamic discharge less SMB to 2100 (Gt); q_2, Q_2 converted to ice flux rate, total dynamic flux less SMB to 2100 (km^3/year); Q_2, total dynamic discharge less 10× SMB to 2100 (Gt); and q_2, Q_2 converted to ice flux rate, total dynamic flux less 10× SMB to 2100 (km^3/year).

<table>
<thead>
<tr>
<th>SLR target</th>
<th>Q (Gt)</th>
<th>q (km^3/year)</th>
<th>Q_2 (Gt)</th>
<th>q_2 (km^3/year)</th>
<th>Q_2 (Gt)</th>
<th>q_2 (km^3/year)</th>
</tr>
</thead>
<tbody>
<tr>
<td>2 m</td>
<td>21.5</td>
<td>724,000</td>
<td>8,650</td>
<td>698,164</td>
<td>8,341</td>
<td>652,464</td>
</tr>
<tr>
<td>5 m</td>
<td>53.8</td>
<td>1,810,000</td>
<td>21,625</td>
<td>1,784,165</td>
<td>21,317</td>
<td>1,738,464</td>
</tr>
</tbody>
</table>
range from 26.8 km/year to 125 km/year, depending on the scenario considered (Table 2 and supporting online material (SOM)). These velocities must be achieved immediately on all outlets considered and held at that level until 2100. Delays in the onset of rapid motion increase the required velocity further (fig. S1).

The scenario velocities far exceed the fastest motion exhibited by any Greenland outlet glacier. For example, the near-doubling of ice discharge from Jakobshavn Glacier in 2004–2005 was associated with an acceleration to 12.6 km/year (7). Similarly, a temporary 80% increase in the speed near the terminus of Kangerdlugssuaq produced a velocity of 14.6 km/year (6). A comparison of calculated (Table 2) and observed (1.23 km/year) average velocities shows that calculated values for a 2-m SLR exceed observations by a factor of 22 when considering all gates and inflated SMB and by a factor of 40 for the marine gates without inflated SMB, which we consider to be the more likely scenario. With the exception of discharge through all gates at inflated SMB (26.8 km/year), none of the velocity magnitudes shown in Table 2 has ever been observed anywhere, even over short time periods. The highest observed velocities have occurred at surging glaciers, including circa (ca.) 70 m/day (25.5 km/year) at Variegated Glacier (17) and 105 m/day (38.3 km/year) at Medvezhy Glacier (18), but were held only for brief periods (hours to days). Although no physical proof is offered that the velocities given in Table 2 cannot be reached or maintained over century time scales, such behavior lies far beyond the range of observations and at the least should not be adopted as a central working hypothesis. Calculations are made only for Greenland because Greenland’s outlet glaciers are well constrained by bed topography, which (despite the uncertainties mentioned) is well known in comparison to much of the Antarctic Ice Sheet and the Antarctic Peninsula and virtually all of the marine-terminating glaciers and ice caps (GIC) exclusive of Greenland and Antarctica. In order to estimate how these constraints influence projections of total SLR to 2100, we calculated a zero-order range of eustatic SLR from all land ice sources. Because marine-grounded channels are not well defined in many other locations, we made approximations and scaling arguments to arrive at a range of values for total eustatic SLR and, including reasonable projections of steric SLR, a range of estimates of total SLR to 2100.

Most of the marine-based ice in West Antarctica is held behind the Ross and Filchner-Ronne ice shelves, which we consider unlikely to be removed by climate or oceanographic processes within the next century [e.g., (19)]. The Amundson Basin coast [including Pine Island Glacier (PIG) and Thwaites Glacier], however, is not confined by large ice shelves and contains about 1.5 m sea level equivalent (5.43 x 1015 Gt) (20). The aggregate cross-sectional area of PIG and Thwaites Glacier is ca. 120 km2 (20). The average velocity in this region is 2 km/year (table S2), higher than the average velocity of all Antarctic ice streams [0.65 km/year (19)]. An average (present day to 2100) velocity of 53.6 km/year is required to discharge 1.5 m sea-level equivalent through the PIG and Thwaites glacier gates by 2100, again far greater than any observed glacier velocity.

We present three scenarios by combining likely projection methods that we believe roughly bracket the range of potential near-future SLR outcomes (SOM). These are not true limiting cases but give a good sense of the potential variability of total SLR due to dynamic discharge effects.

SLR scenario Low 1 represents a low-range estimate based on specific adjustments to dynamic discharge in certain potentially vulnerable areas. We assumed a doubling of outlet glacier velocities in Greenland and PIG/Thwaites within the first decade and no change from present-day discharge values at Lambert/Amery. SMB for Greenland, the Antarctic Peninsula, and GIC was accelerated at present-day rates of SMB change, and, lacking more directly applicable constraints, dynamic discharge for the Antarctic Peninsula and GIC was calculated by scaling dynamic discharge to SMB by using the ratio of 1.31 as computed for Greenland (SOM). The net result, including thermal expansion, is 785 mm by 2100 (Table 3).

A second low-range scenario (Low 2) shows the effect of varying our assumptions; for this, we simply integrated presently observed rates of change forward in time. We calculated Greenland’s contribution as for Low 1 but accelerated the present-day net discharge for Antarctica (East/West Antarctic Peninsula) forward at the present-day rate of change given by (19). The GIC contribution was also calculated by accelerating the present-day net discharge at the current rate of change, with values from (15). The net result, including thermal expansion, is 833 mm by 2100 (Table 3).

SLR scenario High 1 combines all eustatic sources taken at high but reasonable values. No firm highest possible value can be determined for SMB or dynamics; the values chosen represent upper limits of likely behavior on the century time scale. Greenland SMB was accelerated at present-day rates of change, but dynamic discharge was calculated by accelerating outlet glacier velocities by an order of magnitude in the first decade. In Antarctica, PIG/Thwaites was accelerated from present-day net discharge (19) in the first decade and held thereafter to the highest outlet glacier velocity observed anywhere [14.6 km/year (6)], and Lambert/Amery was accelerated from present-day net discharge (19) in the first decade by an order of magnitude and held thereafter. Antarctic Peninsula and GIC were calculated by scaling dynamic discharge at the dynamics-to-SMB ratio computed for Greenland; this ratio is larger (6.42) than in case Low 1 because Greenland’s dynamic discharge is larger. The net result, including thermal expansion, is 2008 mm by 2100 (Table 3).

On the basis of calculations presented here, we suggest that an improved estimate of the range of SLR to 2100 including increased ice dynamics lies between 0.8 and 2.0 m. We emphasize that assumptions made to arrive here contain substantial uncertainties, and many other scenarios and combinations of contributions could be considered. However, the net eustatic SLR from other combinations explored fell within the range given in Table 3. Hence, these values give a context and starting point for refinements in SLR forecasts on the basis of clearly defined assumptions and offer a more plausible range of estimates than those neglecting the dominant ice dynamics term. Certain potentially significant sinks and sources of SLR, such as terrestrial water storage, are still absent altogether. Among the uncertainties explored, the potential for dynamic response from GIC is comparable in magnitude to dynamic response from Greenland or Antarctica but is exceptionally poorly constrained by basic observations. Without better knowledge of the number, size, and catchment areas of marine-based outlet glaciers in the GIC category, improvements on the estimates made here will be very difficult.

### Table 3. SLR projections based on kinematic scenarios. Thermal expansion numbers are from (22).

<table>
<thead>
<tr>
<th>Scenario</th>
<th>Greenland</th>
<th>Antarctica</th>
<th>Glaciers/ice caps</th>
</tr>
</thead>
<tbody>
<tr>
<td>Low 1</td>
<td>93</td>
<td>108</td>
<td>94</td>
</tr>
<tr>
<td>Low 2</td>
<td>93</td>
<td>108</td>
<td>80</td>
</tr>
<tr>
<td>High 1</td>
<td>467</td>
<td>394</td>
<td>471</td>
</tr>
<tr>
<td>Low 2</td>
<td>71</td>
<td>16</td>
<td>80</td>
</tr>
<tr>
<td>High 1</td>
<td>71</td>
<td>12</td>
<td>80</td>
</tr>
<tr>
<td>Low 2</td>
<td>538</td>
<td>129</td>
<td>551</td>
</tr>
<tr>
<td>High 1</td>
<td>165</td>
<td>129</td>
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<tr>
<td>High 1</td>
<td>538</td>
<td>129</td>
<td>551</td>
</tr>
</tbody>
</table>

### References and Notes
Apobec3 Encodes Rfv3, a Gene Influencing Neutralizing Antibody Control of Retrovirus Infection

Mario L. Santiago, Mauricio Montano, Robert Benitez, Ronald J. Messer, Wes Yonemoto, Bruce Chesebro, Kim J. Hasenkrug, Warner C. Greene

Recovery from Friend virus 3 (Rfv3) is a single autosomal gene encoding a resistance trait that influences retroviral neutralizing antibody responses and viremia. Despite extensive research for 30 years, the molecular identity of Rfv3 has remained elusive. Here, we demonstrate that Rfv3 is encoded by Apobec3. Apobec3 maps to the same chromosome region as Rfv3 and has broad inhibitory activity against retroviruses, including HIV. Not only did genetic inactivation of Apobec3 convert Rfv3-resistant mice to a susceptible phenotype, but Apobec3 was also found to be naturally disabled by aberrant messenger RNA splicing in Rfv3-susceptible strains. The link between Apobec3 and neutralizing antibody responses highlights an Apobec3-dependent mechanism of host protection that might extend to HIV and other human retroviral infections.

The study of viral resistance factors has provided important insights into the evolutionary strategies of defense used by mammalian hosts (1–5). Recovery from Friend virus (FV) virus 3 (Rfv3) was first identified as a resistance trait in 1978 (6, 7), and later genetic studies showed that the phenotypes of decreased viremia and FV-specific neutralizing antibody responses segregated as a single gene (8). Because the generation of neutralizing antibodies is critical for recovery from FV infection (1, 9) and a desired but often unrealized outcome in various retroviral infections, including HIV-1, we have focused our efforts on identifying the gene encoding Rfv3. The Rfv3 locus maps to a 0.83-centimorgan region of chromosome 15 (fig. S1A) (10–12), which contains at least 61 annotated genes (table S1), one of which is murine Apobec3 (mA3), a member of a family of deoxycytidine deaminases with antiretroviral and antiretroelement activity [as reviewed in (13)]. This fact, along with the presence of substantial polymorphism in mA3 (table S1), led us to consider mA3 as a prime candidate for Rfv3. Because Rfv3 has no described in vitro phenotype, our investigation required the generation of mA3-deficient mice (14). First, an inactivated mA3 gene (fig. S1B) was introduced into the Rfv3s/s C57BL/6 (B6) background to test its ability to act as a defect Rfv3 allele in matings with Rfv3r/r mice (table 1). Because the Rfv3 resistance trait is dominant over susceptibility (7), Rfv3s/s F1 offspring should control viremia and mount effective neutralizing antibody responses. Conversely, if mA3 encodes Rfv3, then the gene from a mA3−/− parent will be null, and the resultant F1 offspring with an Rfv3−/− genotype are predicted to exhibit higher levels of viremia and weaker neutralizing antibody responses to test this possibility, B6 x BALB/c F1 offspring were infected with FV and plasma viremia levels were measured. At 7 days post infection (dpi), the F1 mice containing an inactivated mA3 gene exhibited levels of viremia 15 times as high as their congenic partners carrying the wild-type mA3 allele (Fig. 1A). These high viral loads in mA3−/− F1 mice were comparable to FV levels found in fully susceptible Rfv3s/s BALB/c parental mice. Thus, mA3 is a restriction factor contributing to the

Table 1. FV infection characteristics of various mouse strains used in this study.

<table>
<thead>
<tr>
<th>Type</th>
<th>Strain</th>
<th>General FV susceptibility</th>
<th>Viremia</th>
<th>Rfv3</th>
<th>Neutralizing antibody</th>
<th>H-2*</th>
<th>Cell-mediated immunity</th>
<th>Fv2‡</th>
<th>Splenomegaly induction</th>
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</thead>
<tbody>
<tr>
<td>Wild-type</td>
<td>C57BL/6 (B6)</td>
<td>Resistant</td>
<td>Resistant</td>
<td>r/r</td>
<td>High</td>
<td>b/b</td>
<td>High</td>
<td>r/r</td>
<td>No</td>
</tr>
<tr>
<td></td>
<td>BALB/c</td>
<td>Susceptible</td>
<td>Chronic</td>
<td>s/s</td>
<td>Low</td>
<td>d/d</td>
<td>Very low</td>
<td>s/s</td>
<td>Yes</td>
</tr>
<tr>
<td></td>
<td>A.BY</td>
<td>Susceptible</td>
<td>Chronic</td>
<td>s/s</td>
<td>Low</td>
<td>b/b</td>
<td>High</td>
<td>s/s</td>
<td>Yes</td>
</tr>
<tr>
<td></td>
<td>129/Ola†</td>
<td>Susceptible</td>
<td>Resistant</td>
<td>r/r</td>
<td>High</td>
<td>b/b</td>
<td>High</td>
<td>r/r</td>
<td>No</td>
</tr>
<tr>
<td>Fv3 hybrids</td>
<td>B6 x BALB/c</td>
<td>Susceptible</td>
<td>Acute</td>
<td>r/s</td>
<td>High</td>
<td>b/d</td>
<td>Low</td>
<td>r/s</td>
<td>Yes</td>
</tr>
<tr>
<td></td>
<td>B6 x A.BY</td>
<td>Susceptible</td>
<td>Acute</td>
<td>r/s</td>
<td>High</td>
<td>b/b</td>
<td>High</td>
<td>r/s</td>
<td>Yes</td>
</tr>
</tbody>
</table>

*H-2 is the murine major histocompatibility complex that dictates cell-mediated immunity against FV (5, 15).
†Fv2 is a dominant FV susceptibility gene that facilitates splenomegaly induction through aberrant signaling in erythroblasts (4).
‡FV susceptibility data on 129/Ola were based on results from this study (fig. 2C) (14). The cell-mediated immune response of this strain was inferred from its H-2 haplotype.

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